

# High Temperature Cuprate Superconductors

## Theoretical Physics Year 4 Project

T. K. Kingsman

School of Physics and Astronomy  
University of Birmingham

March 1, 2015

# Outline

- 1 Introduction
  - Cuprate Structure
  - Phase Diagram
  - Copper Oxide Layer
- 2 Model
  - Hamiltonian
  - Higher Order
  - Improved Hamiltonian
- 3 Electron Doping
  - Hubbard Model
- 4 Hole Doping
  - Determining Limits
  - Unphysical Limit:  $\text{Cu}^{3+}$
  - Physical Limit:  $\text{Cu}^+$

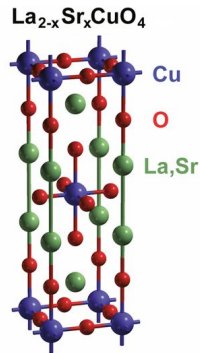
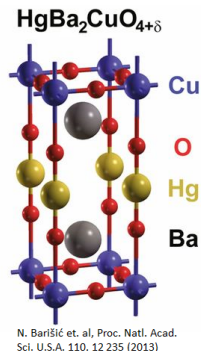
# Outline

- 1 Introduction
  - Cuprate Structure
  - Phase Diagram
  - Copper Oxide Layer
- 2 Model
  - Hamiltonian
  - Higher Order
  - Improved Hamiltonian
- 3 Electron Doping
  - Hubbard Model
- 4 Hole Doping
  - Determining Limits
  - Unphysical Limit:  $\text{Cu}^{3+}$
  - Physical Limit:  $\text{Cu}^+$

# High TC Cuprates - General Structure

Cuprate superconductors have a layered structure.

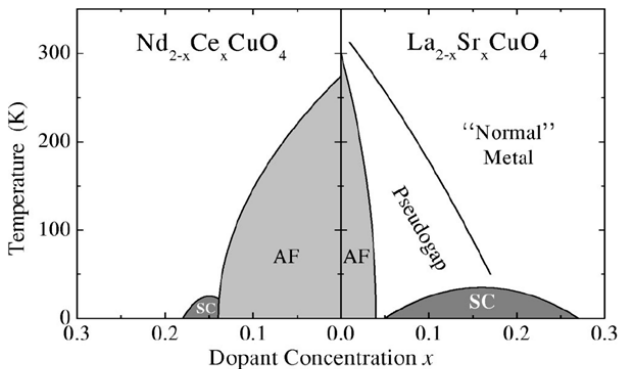
- We introduce extra atoms - **doping**.
- Additional electrons or holes are released.
- A change in the overall charges occurs.
- The copper oxide layers become doped.



The key to the universal properties are the layers of copper oxide.

# Phase Diagram

- Cuprate superconductors can either be **electron** or **hole doped**.
- Experiments show a big difference between the two types.

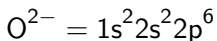
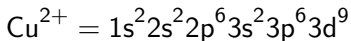


Lee, P., N. Nagaosa, and X.-G. Wen, 2006, Rev. Mod. Phys. 78, 17.

- Why is there an asymmetry in the phase diagram ?

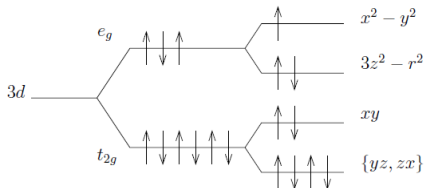
# Relevant Orbitals

Consider the parent compound  $\text{La}_2\text{CuO}_4$  - from  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$   
 Charges on the copper and oxygen ions are,

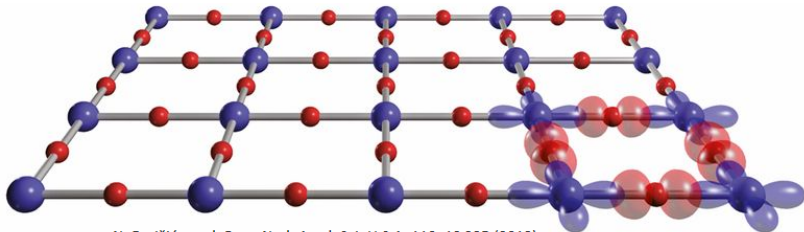


3d and 2p orbitals are further split by  $\text{La}_2\text{CuO}_4$  structure

- Key Cu orbital:  $3d_{x^2-y^2}$
- Key O orbitals:  $2p_x, 2p_y$



# 3 Band Model



N. Barišić et. al, Proc. Natl. Acad. Sci. U.S.A. 110, 12 235 (2013)

**Figure:** A 3 band model of  $\text{CuO}_2$  square plane.

## Key Points

- Every copper site is occupied by a single hole.
- Every oxygen site is unoccupied by holes.
- Doped electrons form  $3d^{10}$  orbitals.
- Doped holes form  $2p^5$  orbitals.

# Outline

- 1 Introduction
  - Cuprate Structure
  - Phase Diagram
  - Copper Oxide Layer
- 2 Model
  - Hamiltonian
  - Higher Order
  - Improved Hamiltonian
- 3 Electron Doping
  - Hubbard Model
- 4 Hole Doping
  - Determining Limits
  - Unphysical Limit:  $\text{Cu}^{3+}$
  - Physical Limit:  $\text{Cu}^+$

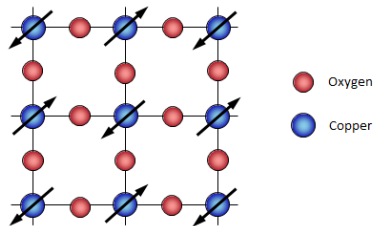


# Hamiltonian

We propose a model for the ground state of the copper oxide layer.

$$\mathcal{H}_0 = -\Delta \sum_{i\sigma} d_{i\sigma}^\dagger d_{i\sigma} + U \sum_i d_{i\sigma}^\dagger d_{i\sigma} d_{i\bar{\sigma}}^\dagger d_{i\bar{\sigma}}$$

- $\Delta$  = the energy needed to place a single hole on a copper site
- $U$  = the energy needed to place a second hole on a copper site



Lee, P., N. Nagaosa, and X.-G. Wen, 2006, Rev. Mod. Phys. 78, 17.

# Higher Order Contributions

- We have a degenerate ground state that we wish to maintain.
- We want to exclude terms of order  $V$  and include order  $V^2$ .

$$\mathcal{H}_0 = -\Delta \sum_{i\sigma} d_{i\sigma}^\dagger d_{i\sigma} + U \sum_i d_{i\sigma}^\dagger d_{i\sigma} d_{i\bar{\sigma}}^\dagger d_{i\bar{\sigma}}$$

$$\mathcal{H}_1 = V \sum_{\langle ij \rangle \sigma} (d_{i\sigma}^\dagger p_{j\sigma} + p_{j\sigma}^\dagger d_{i\sigma})$$

$$\mathcal{H}_2 = ?$$

- How do we obtain a form for  $\mathcal{H}_2$  ?

# Canonical Transformation

This is the canonical transformation, a rotation in Hilbert space.

$$f(\lambda) = e^{\lambda S} (H_0 + H_1) e^{-\lambda S}$$

Evaluating out we get,

$$f(1) = H_0 + (H_1 + [S, H_0]) + ([S, H_1] + \frac{1}{2}[S, [S, H_0]]) + \dots$$

Eliminating the order V dependence gives us,

$$\mathcal{H}_1 + [S, H_0] = 0 \Rightarrow \mathcal{H}_1 = [H_0, S]$$

As a result we get that the new Hamiltonian is,

$$\mathcal{H}_2 = \frac{1}{2}[S, H_1]$$

# Hopping Hamiltonian

The extension to the Hamiltonian contains two separate terms,

$$H_2 = + \frac{V^2}{\Delta} \sum_{\langle im \rangle \sigma} \sum_{\langle ij \rangle} p_{j\sigma}^\dagger p_{m\sigma} - \frac{UV^2}{\Delta(U - \Delta)} \sum_{\langle im \rangle} \sum_{\langle ij \rangle} (d_{i\uparrow}^\dagger p_{j\downarrow}^\dagger - d_{i\downarrow}^\dagger p_{j\uparrow}^\dagger) (p_{m\downarrow} d_{i\uparrow} - p_{m\uparrow} d_{i\downarrow})$$

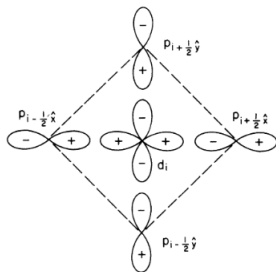
- Term 1 is an **oxygen hopping** term.
- Term 2 is a **singlet hopping** term.

## Note

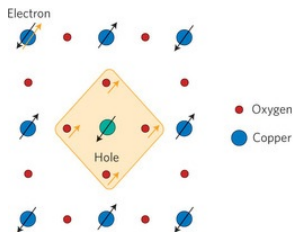
- Valid for both electron and hole doping.

# Physical Interpretation

- Our Hamiltonian also describes a singlet and a triplet solution.
- The lowest energy solution is the **Zhang-Rice Singlet**.



Zhang, F. C. & Rice, T. M.  
Phys. Rev. B 37, 3759–3761 (1988)



J. Orenstein and A. Vishwanath,  
Nat. Phys. 6, 566 (2010)

- A doped hole resonates around an adjacent copper hole.
- Have been observed experimentally.

# Outline

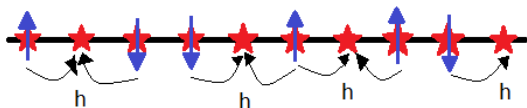
- 1 Introduction
  - Cuprate Structure
  - Phase Diagram
  - Copper Oxide Layer
- 2 Model
  - Hamiltonian
  - Higher Order
  - Improved Hamiltonian
- 3 Electron Doping
  - Hubbard Model
- 4 Hole Doping
  - Determining Limits
  - Unphysical Limit:  $\text{Cu}^{3+}$
  - Physical Limit:  $\text{Cu}^+$

# Hubbard Model Limit

For electron doping, our model reduces to the Hubbard model.

$$\mathcal{H}_2 = -\frac{V^2}{\Delta} \sum_{\langle il \rangle \sigma} d_{l\sigma}^\dagger d_{i\sigma}$$

- Doped electrons will reside on copper sites e.g. removes holes.
- We want to try and solve this model on various geometries.



- Spin order maintained  $\Rightarrow$  concertina effect

$$\mathcal{E}_k = -\frac{2V^2}{\Delta} \cos k$$

# Outline

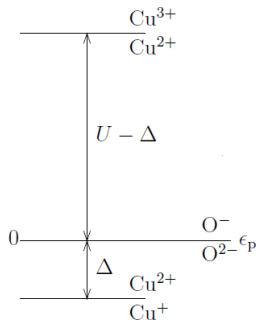
- 1 Introduction
  - Cuprate Structure
  - Phase Diagram
  - Copper Oxide Layer
- 2 Model
  - Hamiltonian
  - Higher Order
  - Improved Hamiltonian
- 3 Electron Doping
  - Hubbard Model
- 4 Hole Doping
  - Determining Limits
  - Unphysical Limit:  $\text{Cu}^{3+}$
  - Physical Limit:  $\text{Cu}^+$



# Limiting Cases

There are two possibilities we need to explore for hole doping:

- Produces  $\text{Cu}^{3+}$  ions:  $\Delta \approx U$ 
  - Easy to double occupy sites.
  - Easier to solve.
  - Unphysical limit.
- Produces  $\text{Cu}^+$  ions:  $\Delta \ll U$ 
  - Never double occupy sites.
  - Harder to solve.
  - Physical limit.



## Line

- In the  $\text{Cu}^{3+}$  limit our Hamiltonian turns into the form of a singlet.

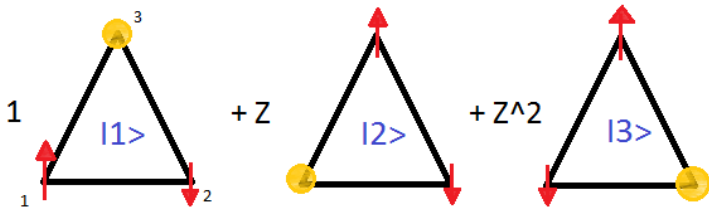
$$\mathcal{H}_2 = -X \sum_{\langle im \rangle} \sum_{\langle ij \rangle} (d_{i\uparrow}^\dagger p_{j\downarrow}^\dagger - d_{i\downarrow}^\dagger p_{j\uparrow}^\dagger) (p_{m\downarrow} d_{i\uparrow} - p_{m\uparrow} d_{i\downarrow})$$

- Solving for the linear chain we get a Hubbard like solution.
- What happens if we solve on more complicated geometries?

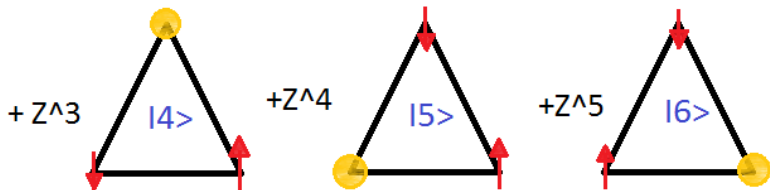
Unphysical Limit:  $\text{Cu}^{3+}$ 

# Triangles

Consider 1 hole and 2 electrons forming the corners of a triangle.

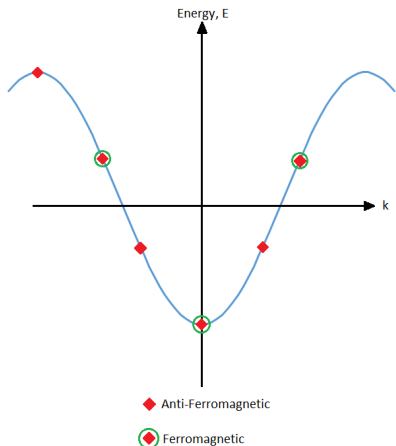


The holes (yellow) move anti-clockwise



# Energy Values - Triangles

Solving for the eigenvalues of the triangular system.



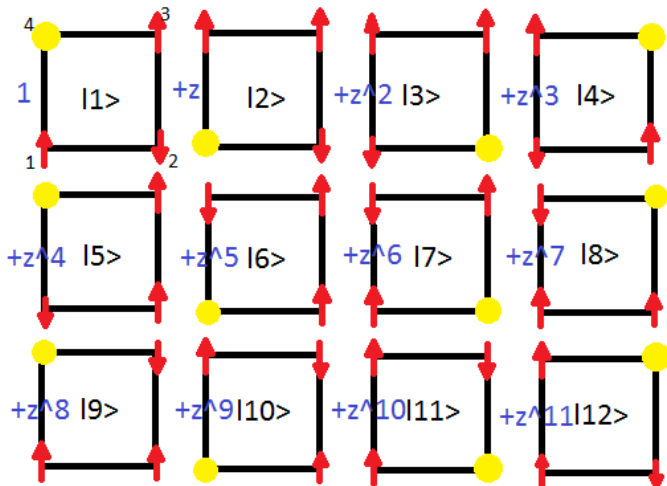
$$\mathcal{E}_k \propto \frac{2V^2}{\Delta} \cos k$$

- Need to combine 2 spin- $\frac{1}{2}$  particles - spin physics
- Either ferromagnetic or anti-ferromagnetic ground state

Unphysical Limit:  $\text{Cu}^{3+}$ 

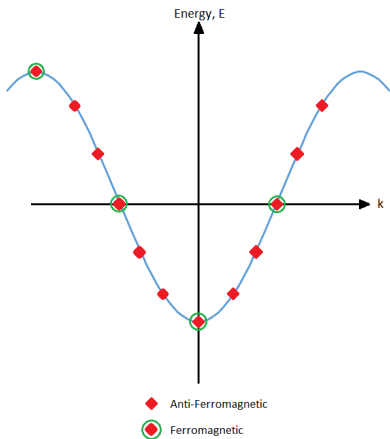
## Squares

Consider 1 hole and 3 electrons forming the corners of a square.



# Energy Values - Squares

Solving for the eigenvalues of the square system.



$$\mathcal{E}_k \propto \frac{2V^2}{\Delta} \cos k$$

- Need to combine 3 spin- $\frac{1}{2}$  particles - spin physics
- Always ferromagnetic ground state
- Differs from the triangle case

# Nagoaka's Theorem

The results we have seen for squares and triangles have been solved in general for even bigger number of sites.

The conditions required are:

- Requires exactly one doped hole
- Sites form a closed loop

## Nagoaka's Theorem - Implications

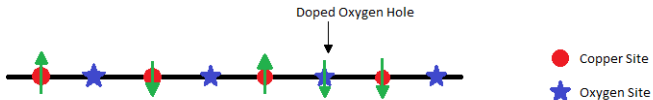
- Even number of sites  $\Rightarrow$  ground state is always ferromagnetic.
- Odd number of sites  $\Rightarrow$  ground state is either of maximum or minimum total spin.

$\text{Cu}^+$  Limit

- In the physical limit we have both hopping and singlet terms.

$$\mathcal{H}_2 = T \sum_{\langle jm \rangle \sigma} p_{j\sigma}^\dagger p_{m\sigma} - T \sum_{\langle im \rangle} \sum_{\langle ij \rangle} (d_{i\uparrow}^\dagger p_{j\downarrow}^\dagger - d_{i\downarrow}^\dagger p_{j\uparrow}^\dagger) (p_{m\downarrow} d_{i\uparrow} - p_{m\uparrow} d_{i\downarrow})$$

- Difficulty is increased as we have two competing interactions.
- We shall illustrate this on a linear chain - surprising result !

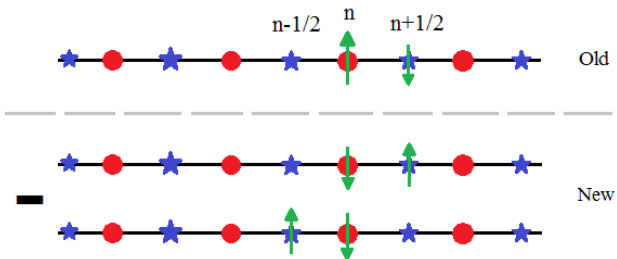




Physical Limit:  $\text{Cu}^+$ 

# Illustrating the Mechanism

- Consider a doped spin which is anti-parallel to its neighbour.



- The holes hop along preserving the spin order - shuffling
- The spin order can be altered - 1st appearance so far !

# Summary

We have proposed a Hamiltonian for our 3 band model.

- Electron doping and  $\text{Cu}^{3+}$  limit both yield Hubbard model.
- $\text{Cu}^+$  limit provides something different.

Our  $\text{Cu}^+$  limit can provide a mechanism for spin exchange.

- Could provide a answer for asymmetry in phase diagram.

Can we find an answer for the linear chain ground state to test experimentally ?