

Part2: Dissipative Systems and Non-Equilibrium Bose-Einstein Condensation

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1 Non-equilibrium condensation in dissipative environment

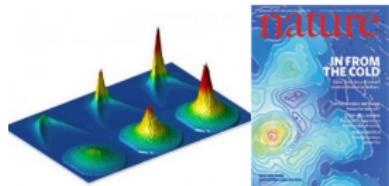
- Polariton model and Keldysh field theory
- Mean-field equations
- Connections of mean-field equations to other limits
- Phase transition and fluctuations
- Superfluidity

Overview: Polaritons

- Experiments

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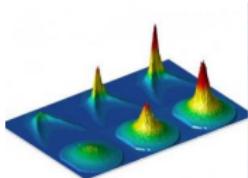
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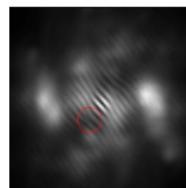
[Kasprzak, et al., Nature
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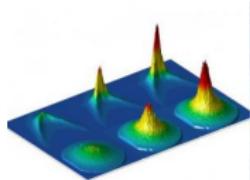
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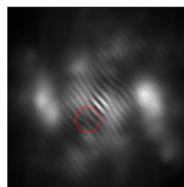
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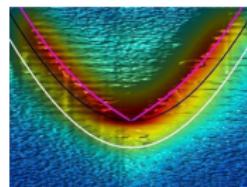
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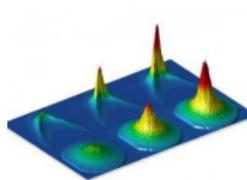
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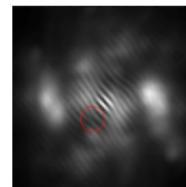
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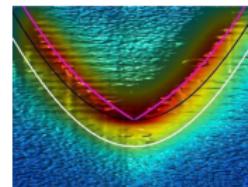
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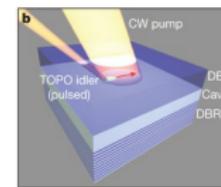
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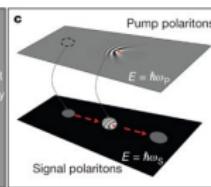
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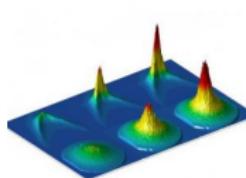


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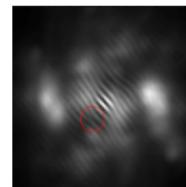


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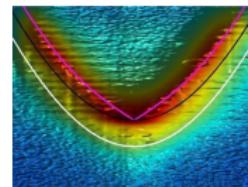
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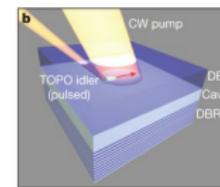
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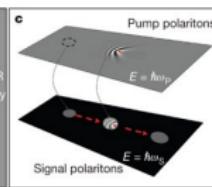
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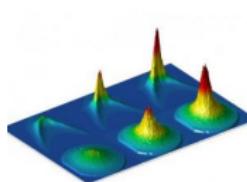
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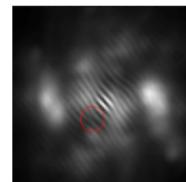
- Theoretically complex

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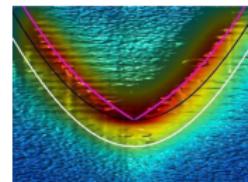
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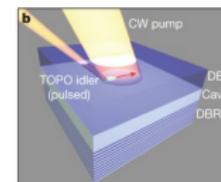
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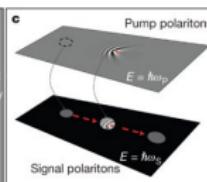
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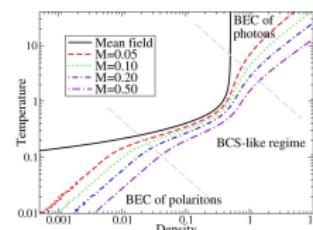
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$E = \hbar\omega_p$

- Theoretically complex

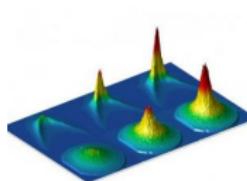
- extended size, internal structure, strong interactions and the **BCS-BEC crossover**



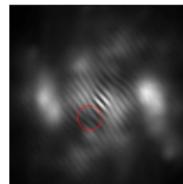
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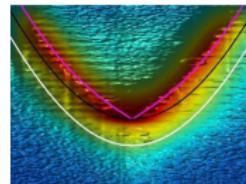
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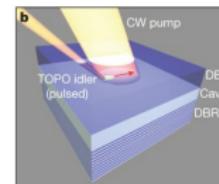
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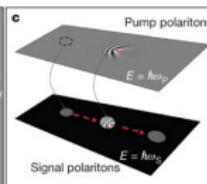
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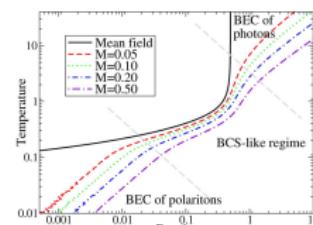
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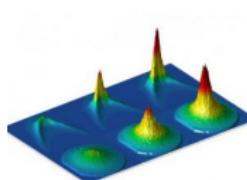
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- 2D vs finite size



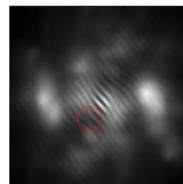
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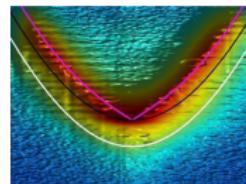
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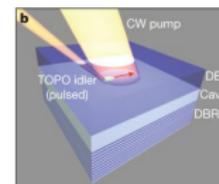
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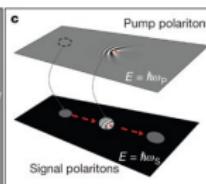
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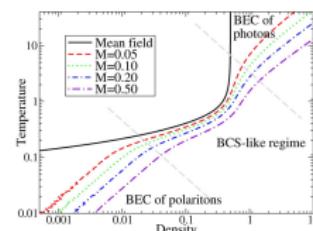
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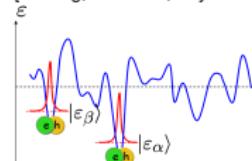
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- extended size, internal structure, strong interactions and the **BCS-BEC crossover**
- 2D vs finite size
- excitonic and photonic disorder



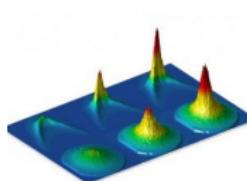
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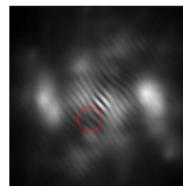
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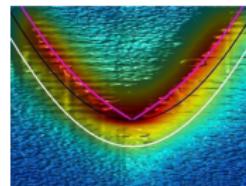
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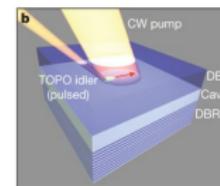
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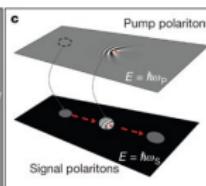
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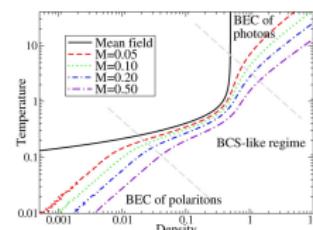
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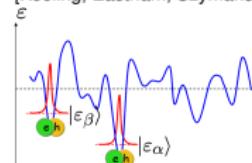
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- extended size, internal structure, strong interactions and the **BCS-BEC crossover**
- 2D vs finite size
- excitonic and photonic disorder
- non-equilibrium and dissipation



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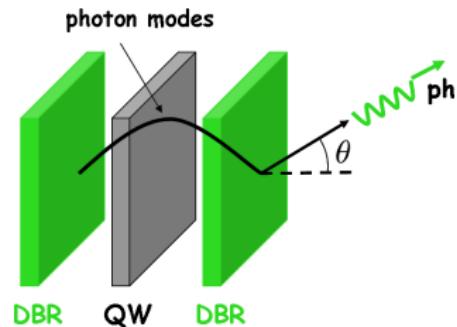
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Non-equilibrium Condensation: model and method

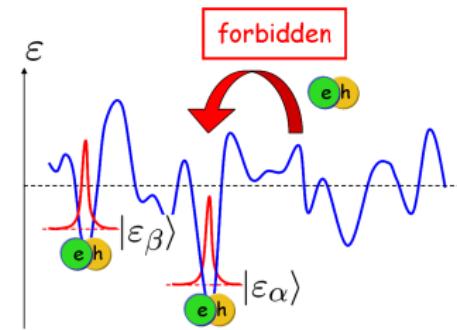
Non-equilibrium Condensation: model and method

Model: $H_{\text{sys}} = \sum_{\mathbf{p}} \omega_{\mathbf{p}} \psi_{\mathbf{p}}^\dagger \psi_{\mathbf{p}}$



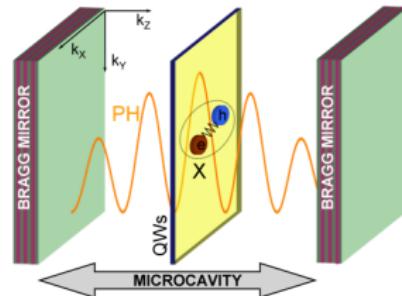
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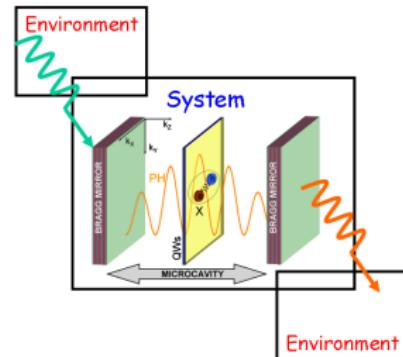
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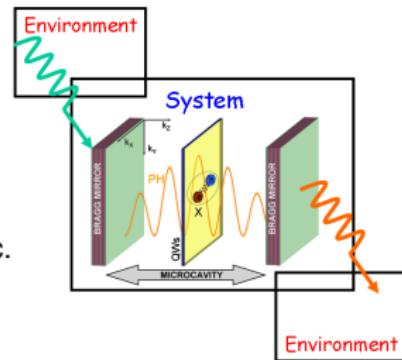
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Schematically: pump γ , decay κ

$$H_{\text{sys,bath}} \simeq \sum_{\mathbf{p}, \mathbf{k}} \sqrt{\kappa} \psi_{\mathbf{p}} \psi_{\mathbf{k}}^\dagger + \sum_{\alpha, \mathbf{k}} \sqrt{\gamma} (a_{\alpha}^\dagger A_{\mathbf{k}} + b_{\alpha}^\dagger B_{\mathbf{k}}) + \text{h.c.}$$



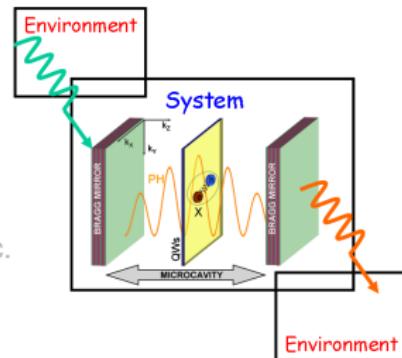
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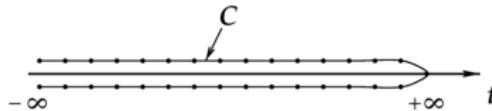
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Method: Keldysh path integral techniques adopted for broken symmetry systems with strong dissipation



Saddle-point of non-equilibrium action: mean-field

Equilibrium BEC: described by Gross-Pitaevskii equation (mean-field equation for the condensate)

$$\left(i\partial_t + \frac{\nabla^2}{2m} - V(r) \right) \psi = U|\psi|^2\psi$$

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Non-equilibrium generalisation of Gross-Pitaevskii in BEC and gap equation in BCS regimes (note: now it is complex)

$$(i\partial_t - \omega_c + i\kappa) \psi = \chi(\psi)\psi$$

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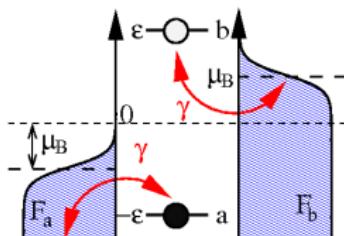
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$$(i\partial_t - \omega_c + i\kappa) \psi = \chi(\psi)\psi$$

For steady-state solution $\psi(\mathbf{r}, t) = \psi e^{-i\mu_S t}$

Susceptibility:

$$\chi(\psi, \mu_S) = -g^2 \gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_S)}{[(\nu - E_\alpha)^2 + \gamma^2][(v + E_\alpha)^2 + \gamma^2]}$$



$$E_\alpha^2 = (\epsilon_\alpha - \frac{1}{2}\mu_S)^2 + g^2|\psi|^2$$

Limits of mean-field equations

Mean-field equations

$$\mu_s - \omega_c + i\kappa = -g^2 \gamma \sum_{\text{excitons}} \int \frac{d\nu}{2\pi} \frac{(F_a + F_b)\nu + (F_b - F_a)(i\gamma + \epsilon_\alpha - \frac{1}{2}\mu_s)}{[(\nu - E_\alpha)^2 + \gamma^2][(v + E_\alpha)^2 + \gamma^2]}$$

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- Large temperature T : laser limit, only imaginary part, gain balances loss

$$\kappa = -g^2 \gamma \sum_{\text{excitons}} \frac{F_b - F_a}{4E_\alpha^2 + 4\gamma^2}$$

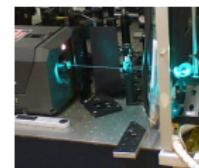
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- Equilibrium limit $\kappa, \gamma \rightarrow 0$:

Limits of mean-field equations

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- Equilibrium limit $\kappa, \gamma \rightarrow 0$: well known gap equation

$$\omega_c - \mu_s = \sum_{\text{excitons}} \frac{g^2}{2E_\alpha} \tanh \left(\frac{\beta E_\alpha}{2} \right)$$



Limits of mean-field equations

Mean-field equation

$$(i\partial_t - \omega_c + i\kappa) \psi = \chi(\psi, \mu_s) \psi$$

- Local density approximation:

$$\left[i\partial_t + i\kappa - \left(V(r) - \frac{\nabla^2}{2m} \right) \right] \psi(r) = \chi(\psi(r, t)) \psi(r, t)$$

Limits of mean-field equations

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$$\left[i\partial_t + i\kappa - \left(V(r) - \frac{\nabla^2}{2m} \right) \right] \psi(r) = \chi(\psi(r, t)) \psi(r, t)$$

At low density the nonlinear, complex susceptibility can be simplified: **Complex Gross-Pitaevskii equation**

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Limits of mean-field equations

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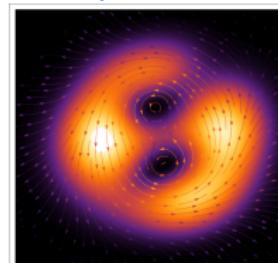
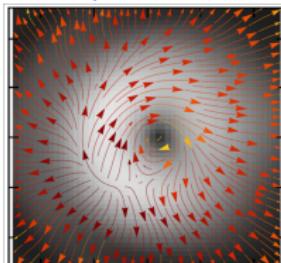
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Vortex lattices, Keeling & Berloff *PRL 2008*

Persistent currents and quantised vortices, D. Sanvitto, F. M. Marchetti, M. H. Szymańska et al, *Nature Physics*, July 2010



Fluctuations: phase transition

Keldysh approach:

$$\mathcal{G}_S = \mathcal{G}_R - \mathcal{G}_A = -i \left\langle [\psi^\dagger, \psi]_- \right\rangle$$

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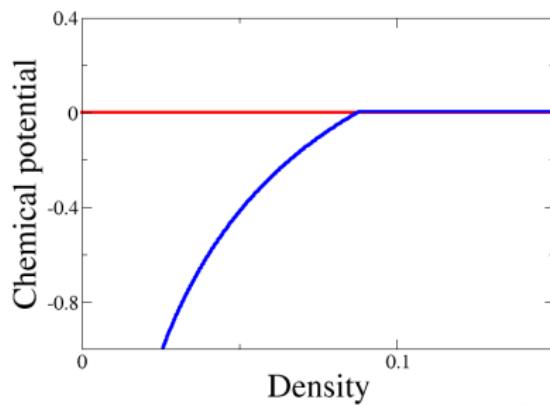
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Equilibrium BEC

$n(\omega) = n_B(\omega)$ i.e $n(\mu)$ diverges

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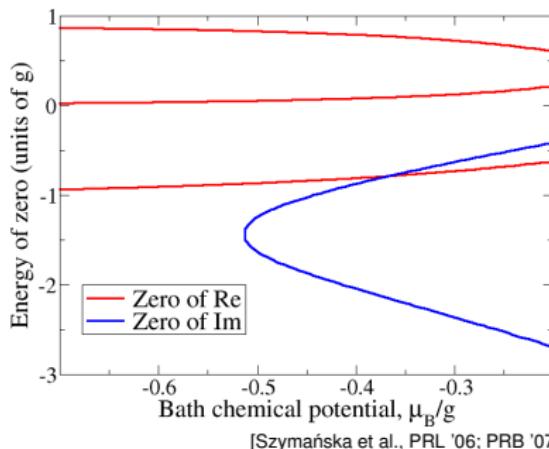
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[Szymańska et al., PRL '06; PRB '07]

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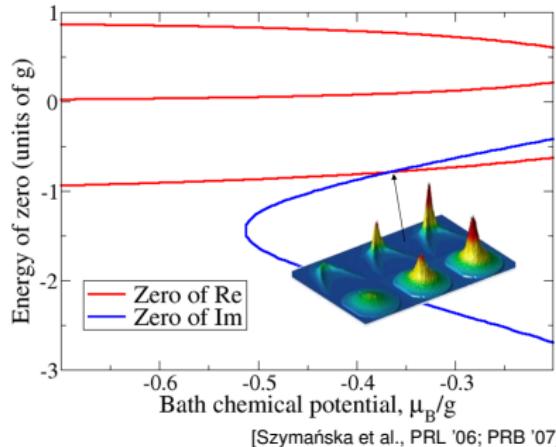
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At transition gap Equation is:

$$\mathcal{G}_R^{-1}(\omega = \mu_S, \mathbf{p} = 0) = 0$$

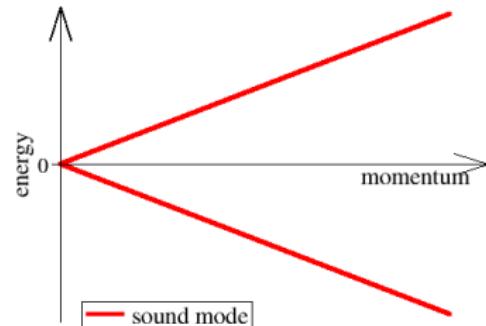
Fluctuations: collective modes and decay of correlations

When condensed

$$\text{Det} \left[\mathcal{G}_R^{-1}(\omega, \mathbf{p}) \right] = \omega^2 - c^2 \mathbf{p}^2$$

Poles:

$$\omega^* = c|\mathbf{p}|$$



[Szymańska et al., PRL '06; PRB '07]

Fluctuations: collective modes and decay of correlations

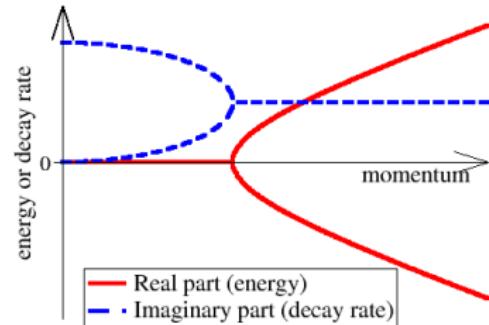
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$$\omega^* = -ix \pm \sqrt{c^2 \mathbf{p}^2 - x^2}$$

Violates Landau criterion!



[Szymańska et al., PRL '06; PRB '07]

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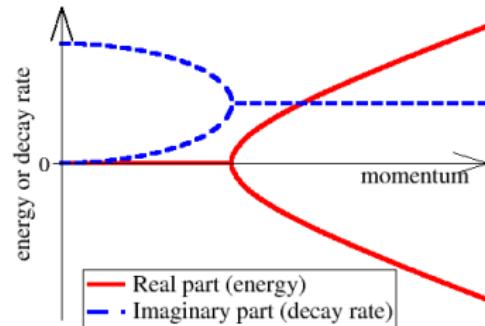
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Correlations (in 2D): $\langle \psi^\dagger(\mathbf{r}, t)\psi(0, r') \rangle \simeq |\psi|^2 \exp[-\mathcal{D}_{\phi\phi}(t, r, r')]$

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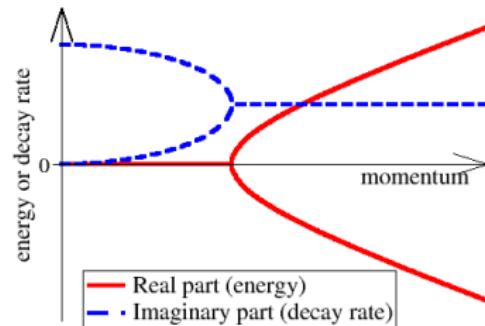
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$$\langle \psi^\dagger(\mathbf{r}, t)\psi(0, 0) \rangle \simeq |\psi|^2 \exp \left[-\eta \begin{cases} \ln(r/\xi) & r \rightarrow \infty, t \simeq 0 \\ \frac{1}{2} \ln(c^2 t / x \xi^2) & r \simeq 0, t \rightarrow \infty \end{cases} \right]$$

$\eta(\text{pump, decay, density})$ and not $\eta(T, \text{density})$ as in equilibrium

[Szymańska et al., PRL '06; PRB '07]

Polariton superfluidity

NATURE|Vol 457|15 January 2009

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Going with the flow

Jonathan Keeling and Natalia G. Berloff

Observations of superfluid behaviour — flow without friction — of unusual character in a condensed-matter system pave the way to investigations of superfluidity in systems that are out of thermal equilibrium.

Table 1 | Superfluidity checklist

	Quantized vortices	Landau critical velocity	Metastable persistent flow	Two-fluid hydrodynamics	Local thermal equilibrium	Solitary waves
Superfluid ^4He /cold atom Bose-Einstein condensate	✓	✓	✓	✓	✓	✓
Non-interacting Bose-Einstein condensate	✓	X	X	X	✓	X
Classical irrotational fluid	X	✓	X	✓	✓	✓
Incoherently pumped polariton condensates	✓	X	?	?	X	?
Parametrically pumped polariton condensates	?	✓	?	?	X	✓

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Conclusion

- Unified description of **condenstes** (BEC and BCS) and **lasers**
- **Bose-Einstein distribution not special** - BEC possible for any “**bosonic**” distribution (i.e divergent at some frequency)

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- Unified description of **condenstes** (BEC and BCS) and **lasers**
- **Bose-Einstein distribution not special** - BEC possible for any “**bosonic**” distribution (i.e divergent at some frequency)
- However, dissipation and driving changes **spectrum of excitations**
 - different collective modes
 - affected correlations i.e spatial and temporal coherence
 - spontaneous vortices
 - changed superfluid properties - **still under experimental investigation**

Extra slides